

TO SUMMARIZE, THE BOTTOM-UP CONSTRUCTION OF THE LEFT/LEFT REQUIRES THE FOLLOWING INGREDIENTS

- 1) DoF $\rightarrow \gamma, g, u, d, s, c, b, e, \mu, \tau, \nu_e, \nu_\mu, \nu_\tau$
- 2) Symmetries $\rightarrow SU(3)_c \otimes U(1)_{em}$
- 3) POWER-COUNTING $\sim m/\sqrt{\Lambda}$ (or m/M_w) WITH LIGHT MASS
 $\sim P/\sqrt{\Lambda}$ (or P/M_w) WITH $P \ll M_w$

ⓐ AT LEADING ORDER IN THE EFT EXPANSION, HOW DOES SUCH LAGRANGIAN LOOK-LIKE?

[THINK $M_w \rightarrow \infty$ LIMIT]

↳ NOTHING BUT QED + QCD LAGRANGIAN, SINCE WEAKLY INTERACTING PARTICLES HAVE BEEN INTEGRATED OUT

ⓑ WILL NEUTRINOS BE PRESENT IN LEFT AT LEADING ORDER?

↳ NO, NEUTRAL UNDER RELEVANT GAUGE SYMMETRY!

$$\mathcal{L}_{\text{LEFT}} = \mathcal{L}_{\text{QED} + \text{QCD}} + \mathcal{O}(v^{-1})$$

↳ EXPANSION PARAMETER

$$\begin{aligned} \mathcal{L}_{\text{QED} + \text{QCD}} &= -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{4} G_{\mu\nu}^A G^{A\mu\nu} \\ &+ \mathcal{O}_{\text{QCD}} \frac{g^2}{32\pi^2} G_{\mu\nu}^A \tilde{G}^{A\mu\nu} + \left[\bar{\psi}_i \not{\partial} \psi \right] \\ &\quad \psi = u, d, e, \nu_i \\ &- \left[\sum_{\psi = u, d, e} \bar{\psi}_R m_\psi \psi_L + \text{h.c.} \right] \end{aligned}$$

WITH THE COVARIANT DERIVATIVE GIVEN BY

$$D_\mu = \partial_\mu + ig T^A G_\mu^A + ie Q A_\mu$$

IF ACCIDENTAL SYMMETRIES OF SM DISREGARDED,
THERE IS A MAJORANA MASS TERM FOR NEUTRINOS WITH $\Delta L=2$

TO CONSTRUCT FULL LEFT, RECIPE IS TO ADD
ALL HIGHER-DIMENSIONAL OPERATORS CONSISTENT WITH
PARTICLE CONTENT AND FUNDAMENTAL SYMMETRIES

$$\mathcal{L}_{\text{LEFT}} = \mathcal{L}_{QED+QCD} + \left[\sum_{d \geq 5} \sum_{i=1}^{n_d} L_i^{(d)} O_i^{(d)} \right]$$

$$L_i^{(d)} \propto V^{4-d} \quad [\text{WILSON COEFFICIENT}]$$

Q IS LEFT AT $O(V^0)$ A CHIRAL THEORY

$\mathcal{L}_{QED+QCD}$ IS VECTOR GAUGE THEORY

→ SAME INTERACTIONS FOR LEFT AND RIGHT-HANDED FERMIONS

$\mathcal{L}_{\text{LEFT}}$ → DIFFERENT INTERACTIONS FOR
LEFT AND RIGHT-HANDED FERMIONS

NOW, LET'S SEE WHICH KIND OF HIGHER-DIMENSIONAL
OPERATORS WE CAN CONSTRUCT IN LEFT

$$X = F_{\mu\nu}, \text{ or } G_{\mu\nu}$$

4 → FERMIONS

Q WHAT HAPPENS AT DIM-5?
WHICH KIND OF OPERATOR DO ARISE?

$$\dim(x) = 2 \quad \dim(\psi) = 3/2$$

$$\dim(O) = 5 \rightarrow O = \psi^2 X$$

$$\dim(O) = 6 \quad O = X^3, \psi^4$$

↳ FOUR-FERMION OPERATORS
ARISING IN FERMION THEORY

Q BASED ON POWER COUNTING, ARE OTHER
TYPES OF OPERATORS POSSIBLE?

YES, INCLUDING DERIVATIVES

$$[\partial] = 1$$

$$\dim(O) = 6 \quad O \sim (\partial X)^2$$

Q WHY THIS IS NOT CONSIDERED?

BECAUSE IT IS REDUNDANT, AND CAN BE REDUCED
TO PREVIOUS LIST OF OPERATORS USING EoM

$$\partial_\mu F^{\mu\nu} = e \int_{\psi=u,d,e} \bar{\psi} \gamma^\nu \psi$$

$$(\partial_\mu G^{\mu\nu})^A = g \int_{\psi=u,d} \bar{\psi} T^A \gamma^\nu \psi$$

WHICH IMPLIES THAT WE CAN REWRITE OPERATORS AS

$$O \sim (\partial X)^2 \sim (\psi^2)^2 \sim \psi^4$$

WHICH IS ALREADY INCLUDED IN THE
LIST

HENCE, WHICH OPERATORS ENTER THE LEFT?

$$\dim(O) = 5 \rightarrow O \sim \psi^2 X$$

$$O_{UG} = \bar{U}_L G^{\mu\nu} T^A U_R G_{\mu\nu}^A$$

$$G_{\mu\nu}^A = \partial_\mu G_\nu^A - \partial_\nu G_\mu^A - g f^{ABC} G_\mu^B G_\nu^C$$

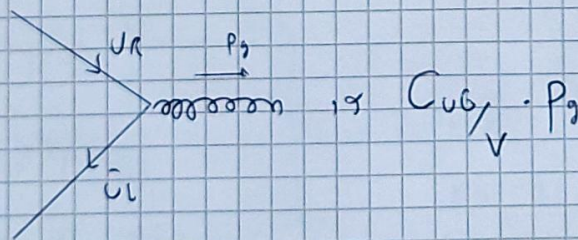
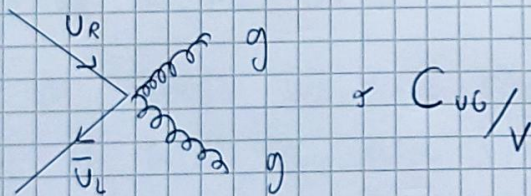
Q) WHERE CAN WE SEE IN THIS OPERATOR THAT THE "FUNDAMENTAL THEORY" IS CHIRAL?

↳ INTERACTIONS DEPEND ON QUARK CHIRALITY

Q) WHAT KIND OF FEYNMAN RULES WOULD THIS OPERATOR IMPLY? CAN WE TRACE THEM BACK TO PROCESSES IN THE "FUNDAMENTAL THEORY"?

$$O_{UG} \sim \psi^2 X \quad X \sim \partial G, G^2$$

$$\sim \psi^2 \partial G, \psi^2 G^2$$



Q) WHAT "SM" PROCESSES ARE RESPONSIBLE?
DISCUSS EXERCISE SESSION

Q) SHOULD WE REPRODUCE IN ALL CASES SM PROCESSES?

$$\dim(O) = G \rightarrow O \sim X^3, \psi^4$$

$$O_G \sim \int^{ABC} G_M^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$$

Q WHAT KIND OF INTERACTIONS?

$$G_\nu^{A\rho} \sim \partial_\mu G, G^2$$

$$O_G \sim (\partial G)^3, (\partial G)^2 G^2, (\partial G) G^4, G^6$$



$$\sim G^3$$

(MULTIJET PRODUCTION @ LHC)

Q DOES THIS FEYNMAN RULE ARISE IF I MATCH THE LEFT TO THE SM?

NO! WE ARE NOT BUILDING THE LEFT BOTTOM-UP SO SOME OPERATORS WILL BE ABSENT IN THE SM

AND OF COURSE WE HAVE FOUR-FERMION OPERATORS, SOME OF THEM WE HAVE SEEN ALREADY IN THE DISCUSSION OF FERMI THEORY

$$(\bar{L}\bar{L})(\bar{R}R)$$

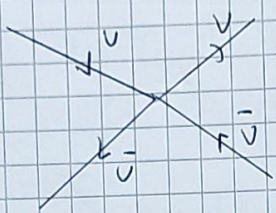
$$(\bar{R}R)(\bar{R}R)$$

$$(\bar{L}L)(\bar{R}R)$$

→ PURELY LEPTONIC,
SEMI-LEPTONIC
NON-LEPTONIC

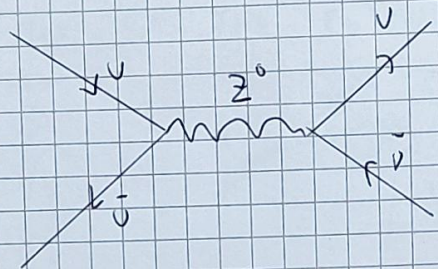
$$O_{\nu\nu}^{V,LL} \propto \frac{v_{HLL}}{\Lambda^2} (\bar{\Psi}_\nu^L \gamma^\mu \Psi_\nu^L) (\bar{\Psi}_\nu^L \gamma_\mu \Psi_\nu^L)$$

COUPLED NEUTRINOS WITH UP-TYPE QUARKS



$$\propto C_{\nu u}^{V,LL} / \sqrt{2}$$

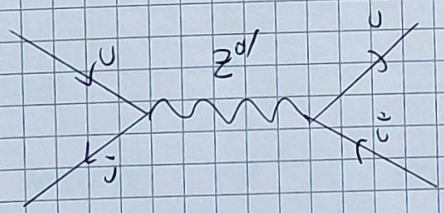
Q DOES THIS OPERATOR ARISE FROM A SM REACTION? IF SO, WHICH ONE?



HENCE IN THE MATCHING TO SM WE FIND NON-ZERO VALUE OF COEFFICIENT WILSON

$$C_{\nu u}^{V,LL} / \sqrt{2} \propto \frac{g_{EW}^2}{M_Z^2} \neq 0$$

THIS WILSON COEFFICIENT CAN ALSO RECEIVE BSM CONTRIBUTIONS EG OF THE FORM



$$m_{Z'} \gg m_Z$$

$$\frac{C_{\nu u}^{V,LL}}{\sqrt{2}} \propto \frac{g_{EW}^2}{M_Z^2} + \alpha \frac{g_{Z'}^2}{M_{Z'}^2} + \dots$$

THE LEFT CAN ALSO BE MATCHED TO SMEFT, TO BE DISCUSSED IN SUBSEQUENT LECTURES